

Lösungen des Übungsblattes 6 zur Vorlesung Theoretische Chemie I WS 2018/19 – Übungsblatt 6

1.

$$\varphi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right) \quad , \quad n \in \mathbb{N} \quad (1)$$

$$E_n = \frac{\hbar^2 \pi^2 n^2}{2mL^2} \quad , \quad n \in \mathbb{N} \quad (2)$$

a)

$$\psi_n(x, t) = \varphi_n(x) e^{-i \frac{E_n}{\hbar} t}$$

Die zeitabhängige Schrödingergleichung für Teilchen in einer Box lautet:

$$\hat{H}\psi_n(x, t) = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi_n(x, t) + V(x)\psi_n(x, t) = i\hbar \frac{\partial}{\partial t} \psi_n(x, t) \quad (3)$$

Innerhalb der Box ist $V(x) = 0$.

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi_n(x, t) = i\hbar \frac{\partial}{\partial t} \psi_n(x, t) \quad (4)$$

Beide Seiten der Gleichung getrennt lösen,

$$\begin{aligned} \text{LHS} &= -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \varphi_n(x) e^{-i \frac{E_n}{\hbar} t} \\ &= -\frac{\hbar^2}{2m} e^{-i \frac{E_n}{\hbar} t} \frac{\partial^2}{\partial x^2} \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right) \\ &= -\frac{\hbar^2}{2m} e^{-i \frac{E_n}{\hbar} t} \cdot \sqrt{\frac{2}{L}} \cdot \frac{\partial}{\partial x} \cos\left(\frac{n\pi x}{L}\right) \cdot \frac{n\pi}{L} \\ &= -\frac{\hbar^2}{2m} e^{-i \frac{E_n}{\hbar} t} \cdot \sqrt{\frac{2}{L}} \cdot \frac{n^2 \pi^2}{L^2} \left(-\sin\left(\frac{n\pi x}{L}\right)\right) \\ &= \frac{\hbar^2}{2m} \frac{n^2 \pi^2}{L^2} \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right) e^{-i \frac{E_n}{\hbar} t} \\ &= \frac{\hbar^2 \pi^2 n^2}{2mL^2} \varphi_n(x) e^{-i \frac{E_n}{\hbar} t} \\ &= E_n \psi_n(x, t) \end{aligned}$$

$$\begin{aligned} \text{RHS} &= i\hbar \frac{\partial}{\partial t} \varphi_n(x) e^{-i \frac{E_n}{\hbar} t} = i\hbar \varphi_n(x) \frac{\partial}{\partial t} e^{-i \frac{E_n}{\hbar} t} \\ &= i\hbar \varphi_n(x) \cdot \left(-\frac{i}{\hbar}\right) E_n e^{-i \frac{E_n}{\hbar} t} \\ &= \varphi_n(x) E_n e^{-i \frac{E_n}{\hbar} t} = E_n \varphi_n(x) e^{-i \frac{E_n}{\hbar} t} \\ &= E_n \psi_n(x, t) \end{aligned}$$

$\therefore$ LHS = RHS. Daher erfüllt $\psi_n(x, t)$ die zeitabhängig Schrödingergleichung des Teilchens im Kasten.

b)

$$\begin{aligned}
\Psi(x, t) &= c_1 \psi_1(x, t) + c_2 \psi_2(x, t) \\
&= c_1 \varphi_1(x) e^{-i \frac{E_1}{\hbar} t} + c_2 \varphi_2(x) e^{-i \frac{E_2}{\hbar} t} \\
&= c_1 \sqrt{\frac{2}{L}} \sin\left(\frac{\pi x}{L}\right) e^{-i \frac{E_1}{\hbar} t} + c_2 \sqrt{\frac{2}{L}} \sin\left(\frac{2\pi x}{L}\right) e^{-i \frac{E_2}{\hbar} t}
\end{aligned} \tag{5}$$

$$\begin{aligned}
|\Psi(x, t)|^2 &= \Psi^*(x, t) \Psi(x, t) \\
&= \left[c_1 \varphi_1(x) e^{-i \frac{E_1}{\hbar} t} + c_2 \varphi_2(x) e^{-i \frac{E_2}{\hbar} t} \right]^* \left[c_1 \varphi_1(x) e^{-i \frac{E_1}{\hbar} t} + c_2 \varphi_2(x) e^{-i \frac{E_2}{\hbar} t} \right] \\
&= \left[c_1^* \varphi_1^*(x) e^{i \frac{E_1}{\hbar} t} + c_2^* \varphi_2^*(x) e^{i \frac{E_2}{\hbar} t} \right] \left[c_1 \varphi_1(x) e^{-i \frac{E_1}{\hbar} t} + c_2 \varphi_2(x) e^{-i \frac{E_2}{\hbar} t} \right] \\
&= c_1^* \varphi_1^*(x) e^{i \frac{E_1}{\hbar} t} c_1 \varphi_1(x) e^{-i \frac{E_1}{\hbar} t} + c_1^* \varphi_1^*(x) e^{i \frac{E_1}{\hbar} t} c_2 \varphi_2(x) e^{-i \frac{E_2}{\hbar} t} \\
&\quad + c_2^* \varphi_2^*(x) e^{i \frac{E_2}{\hbar} t} c_1 \varphi_1(x) e^{-i \frac{E_1}{\hbar} t} + c_2^* \varphi_2^*(x) e^{i \frac{E_2}{\hbar} t} c_2 \varphi_2(x) e^{-i \frac{E_2}{\hbar} t} \\
&= |c_1|^2 |\varphi_1(x)|^2 + |c_2|^2 |\varphi_2(x)|^2 + c_1^* c_2 \varphi_1^*(x) \varphi_2(x) e^{i \left(\frac{E_1 - E_2}{\hbar} \right) t} \\
&\quad + c_2^* c_1 \varphi_2^*(x) \varphi_1(x) e^{i \left(\frac{E_2 - E_1}{\hbar} \right) t}
\end{aligned} \tag{6}$$

$$c_1^* = c_1 \text{ und } c_2^* = c_2 \quad \because c_1, c_2 \in \mathbb{R}$$

$$\begin{aligned}
|\Psi(x, t)|^2 &= c_1^2 \varphi_1^2(x) + c_2^2 \varphi_2^2(x) + c_1 c_2 \varphi_1(x) \varphi_2(x) \left[e^{i \left(\frac{E_2 - E_1}{\hbar} \right) t} + e^{-i \left(\frac{E_2 - E_1}{\hbar} \right) t} \right] \\
&= c_1^2 \varphi_1^2(x) + c_2^2 \varphi_2^2(x) + 2 c_1 c_2 \varphi_1(x) \varphi_2(x) \cos \left[\left(\frac{E_2 - E_1}{\hbar} \right) t \right]
\end{aligned} \tag{7}$$

Die Aufenthaltswahrscheinlichkeitsdichte ist nicht zeitlich konstant. Sie schwingt mit der Frequenz, die der Differenz zwischen den Energieniveaus $\frac{E_2 - E_1}{\hbar}$ entspricht. Im Falle $c_1 = c_2 = \frac{1}{\sqrt{2}}$,

$$\begin{aligned}
\Psi(x, t) &= c_1 \sqrt{\frac{2}{L}} \sin\left(\frac{\pi x}{L}\right) e^{-i \frac{E_1}{\hbar} t} + c_2 \sqrt{\frac{2}{L}} \sin\left(\frac{2\pi x}{L}\right) e^{-i \frac{E_2}{\hbar} t} \\
&= \frac{1}{\sqrt{2}} \sqrt{\frac{2}{L}} \sin\left(\frac{\pi x}{L}\right) e^{-i \frac{E_1}{\hbar} t} + \frac{1}{\sqrt{2}} \sqrt{\frac{2}{L}} \sin\left(\frac{2\pi x}{L}\right) e^{-i \frac{E_2}{\hbar} t} \\
&= \frac{1}{\sqrt{L}} \sin\left(\frac{\pi x}{L}\right) e^{-i \frac{E_1}{\hbar} t} + \frac{1}{\sqrt{L}} \sin\left(\frac{2\pi x}{L}\right) e^{-i \frac{E_2}{\hbar} t} \\
|\Psi(x, t)|^2 &= \frac{1}{2} \varphi_1^2(x) + \frac{1}{2} \varphi_2^2(x) + 2 \cdot \frac{1}{\sqrt{2}} \cdot \frac{1}{\sqrt{2}} \varphi_1(x) \varphi_2(x) \cos \left[\left(\frac{E_2 - E_1}{\hbar} \right) t \right] \\
&= \frac{1}{2} \varphi_1^2(x) + \frac{1}{2} \varphi_2^2(x) + \varphi_1(x) \varphi_2(x) \cos \left[\left(\frac{E_2 - E_1}{\hbar} \right) t \right]
\end{aligned}$$

c)

$$\langle \hat{x} \rangle = \int_0^L \chi^*(x, t) \hat{x} \chi(x, t) dx \tag{8}$$

Für die erste Eigenfunktion, $\chi(x, t) = \psi_1(x, t)$,

$$\begin{aligned}
\langle \hat{x} \rangle &= \int_0^L \varphi_1^*(x) e^{i \frac{E_1}{\hbar} t} x \varphi_1(x) e^{-i \frac{E_1}{\hbar} t} dx = \int_0^L x \varphi_1^2(x) dx = \frac{2}{L} \int_0^L x \sin^2 \left(\frac{\pi x}{L} \right) dx \\
&= \frac{2}{L} \left[\frac{2L^2 \cdot \frac{\pi^2}{L^2} - 2L \cdot \frac{\pi}{L} \sin \left(2L \cdot \frac{\pi}{L} \right) - \cos \left(2L \cdot \frac{\pi}{L} \right) + 1}{8 \frac{\pi^2}{L^2}} \right] \\
&= \frac{2}{L} \left[\frac{2\pi^2 - 2\pi \sin 2\pi - \cos 2\pi + 1}{\frac{8\pi^2}{L^2}} \right] = \frac{2}{L} \cdot \frac{L^2}{8\pi^2} [2\pi^2 - 0 - 1 + 1] \\
&= \frac{L}{4\pi^2} \cdot 2\pi^2 = \frac{L}{2}
\end{aligned} \tag{9}$$

Für die Superposition aus Aufgabenteil b), $\chi(x, t) = \Psi(x, t)$:

$$\begin{aligned}
\langle \hat{x} \rangle &= \int_0^L \Psi^*(x, t) \hat{x} \Psi(x, t) dx = \int_0^L x |\Psi(x, t)|^2 dx \\
&= \int_0^L x \left\{ c_1^2 \varphi_1^2(x) + c_2^2 \varphi_2^2(x) + 2c_1 c_2 \varphi_1(x) \varphi_2(x) \cos \left[\left(\frac{E_2 - E_1}{\hbar} \right) t \right] \right\} dx \\
&= \underbrace{c_1^2 \int_0^L x \varphi_1^2(x) dx}_I + \underbrace{c_2^2 \int_0^L x \varphi_2^2(x) dx}_{II} + 2c_1 c_2 \cos \left[\left(\frac{E_2 - E_1}{\hbar} \right) t \right] \underbrace{\int_0^L x \varphi_1(x) \varphi_2(x) dx}_{III}
\end{aligned} \tag{10}$$

Die Begriffe getrennt lösen,

$$I = \frac{L}{2} \quad \text{aus Gl.(9)}$$

$$\begin{aligned}
II &= \frac{2}{L} \int_0^L x \sin^2 \left(\frac{2\pi x}{L} \right) dx = \frac{2}{L} \left[\frac{2L^2 \cdot \frac{4\pi^2}{L^2} - 2L \cdot \frac{2\pi}{L} \sin \left(2L \cdot \frac{2\pi}{L} \right) - \cos \left(2L \cdot \frac{2\pi}{L} \right) + 1}{8 \cdot \frac{4\pi^2}{L^2}} \right] \\
&= \frac{2}{L} \left[\frac{8\pi^2 - 4\pi \sin 4\pi - \cos 4\pi + 1}{\frac{32\pi^2}{L^2}} \right] = \frac{2}{L} \cdot \frac{L^2}{32\pi^2} [8\pi^2 - 0 - 1 + 1] \\
&= \frac{L}{16\pi^2} \cdot 8\pi^2 = \frac{L}{2}
\end{aligned}$$

$$\begin{aligned}
\text{III} &= \int_0^L x \sqrt{\frac{2}{L}} \sin\left(\frac{\pi x}{L}\right) \sqrt{\frac{2}{L}} \sin\left(\frac{2\pi x}{L}\right) dx = \frac{2}{L} \int_0^L x \sin\left(\frac{\pi x}{L}\right) \sin\left(\frac{2\pi x}{L}\right) dx \\
&= \frac{2}{L} \cdot \frac{1}{2} \int_0^L x \left[\cos\left(\frac{\pi x}{L} - \frac{2\pi x}{L}\right) - \cos\left(\frac{\pi x}{L} + \frac{2\pi x}{L}\right) \right] dx \\
&= \frac{1}{L} \int_0^L x \left[\cos\left(-\frac{\pi x}{L}\right) - \cos\left(\frac{3\pi x}{L}\right) \right] dx \\
&= \frac{1}{L} \int_0^L x \cos\left(\frac{\pi x}{L}\right) dx - \frac{1}{L} \int_0^L x \cos\left(\frac{3\pi x}{L}\right) dx \\
&= \frac{1}{L} \left[\frac{x \sin\left(\frac{\pi x}{L}\right)}{\frac{\pi}{L}} + \frac{\cos\left(\frac{\pi x}{L}\right)}{\left(\frac{\pi}{L}\right)^2} \right]_0^L - \frac{1}{L} \left[\frac{x \sin\left(\frac{3\pi x}{L}\right)}{\frac{3\pi}{L}} + \frac{\cos\left(\frac{3\pi x}{L}\right)}{\left(\frac{3\pi}{L}\right)^2} \right]_0^L \\
&= \frac{1}{L} \left[\frac{L \sin \pi}{\frac{\pi}{L}} + \frac{\cos \pi}{\frac{\pi^2}{L^2}} - 0 - \frac{\cos 0}{\frac{\pi^2}{L^2}} \right] - \frac{1}{L} \left[\frac{L \sin 3\pi}{\frac{3\pi}{L}} + \frac{\cos 3\pi}{\frac{9\pi^2}{L^2}} - 0 - \frac{\cos 0}{\frac{9\pi^2}{L^2}} \right] \\
&= \frac{1}{L} \left[0 - \frac{1}{\frac{\pi^2}{L^2}} - 0 - \frac{1}{\frac{\pi^2}{L^2}} \right] - \frac{1}{L} \left[0 - \frac{1}{\frac{9\pi^2}{L^2}} - 0 - \frac{1}{\frac{9\pi^2}{L^2}} \right] \\
&= \frac{1}{L} \cdot \frac{L^2}{\pi^2} \cdot (-2) - \frac{1}{L} \cdot \frac{L^2}{9\pi^2} \cdot (-2) = -\frac{2L}{\pi^2} + \frac{2L}{9\pi^2} \\
&= \frac{-18L + 2L}{9\pi^2} = -\frac{16L}{9\pi^2}
\end{aligned}$$

$$\begin{aligned}
\therefore \langle \hat{x} \rangle &= c_1^2 \frac{L}{2} + c_2^2 \frac{L}{2} + 2c_1 c_2 \cos \left[\left(\frac{E_2 - E_1}{\hbar} \right) t \right] \left(-\frac{16L}{9\pi^2} \right) \\
&= \frac{L}{2} (c_1^2 + c_2^2) - \frac{32L}{9\pi^2} c_1 c_2 \cos \left[\left(\frac{E_2 - E_1}{\hbar} \right) t \right] \\
&= \frac{L}{2} - \frac{32L}{9\pi^2} c_1 c_2 \cos \left[\left(\frac{E_2 - E_1}{\hbar} \right) t \right]
\end{aligned} \tag{11}$$

Im Falle $c_1 = c_2 = \frac{1}{\sqrt{2}}$,

$$\begin{aligned}
\langle \hat{x} \rangle &= \frac{1}{2} \cdot \frac{L}{2} + \frac{1}{2} \cdot \frac{L}{2} + 2 \cdot \frac{1}{\sqrt{2}} \cdot \frac{1}{\sqrt{2}} \cos \left[\left(\frac{E_2 - E_1}{\hbar} \right) t \right] \left(-\frac{16L}{9\pi^2} \right) \\
&= \frac{L}{4} + \frac{L}{4} - \frac{16L}{9\pi^2} \cos \left[\left(\frac{E_2 - E_1}{\hbar} \right) t \right] \\
&= \frac{L}{2} - \frac{16L}{9\pi^2} \cos \left[\left(\frac{E_2 - E_1}{\hbar} \right) t \right]
\end{aligned} \tag{12}$$